Let us set:

$$
\overrightarrow{O M}=\sqrt{ } N \mathbf{W}
$$

$\overrightarrow{M C}$ : vector normal to the plane $\left(\mathbf{V}_{1}, \mathbf{V}_{2}\right)$.
By Pythagoras's theorem, we have:

$$
X_{m}^{2}=|\overrightarrow{O C}|^{2}=N|\mathbf{W}|^{2}-|\overrightarrow{M C}|^{2}=N-|\overrightarrow{M C}|^{2}
$$

By elementary geometry and equation (3) we have:

$$
|\overrightarrow{M C}|^{2}=\left(\frac{\text { volume of }\left(\mathbf{V}_{1}, \mathbf{V}_{2}, V N \mathbf{W}\right)}{\text { surface }} \frac{\operatorname{of}\left(\mathbf{V}_{1}, \mathbf{V}_{2}\right)}{)^{2}}=N \frac{\Delta_{3}}{D_{2}} .(B .7)\right.
$$

By a straightforward generalization, we have:

$$
\begin{equation*}
|\overrightarrow{M C}|^{2}=N \frac{\Delta_{m+1}}{D_{m}} \tag{B.7a}
\end{equation*}
$$

Therefore:
$Q_{m}=N \frac{D_{m}-\Delta_{m+1}}{D_{m}}=N-N \frac{\Delta_{m+1}}{D_{m}}=N-|M C|^{2}=X_{m}^{2}$.
We conclude:

$$
\begin{equation*}
Q_{m}=X_{m}^{2} \tag{B.8}
\end{equation*}
$$

is the square of the modulus of the projection of $\mathbf{W}$ upon the 'sub-space' $\left(\mathbf{V}_{1} \ldots \mathbf{V}_{m}\right)$.
(D) Proof of inequality (15a)

Clearly

$$
|\overrightarrow{M C}|^{2} \leq N|\mathbf{W}|^{2}=N
$$

Therefore, from $(B .7 a)$, we obtain

$$
\begin{equation*}
\frac{\Delta_{m+1}}{D_{m}} \leq 1 \tag{B.9}
\end{equation*}
$$

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# The Strengthening of Direct Methods of Crystal Structure Determination by use of Data from Isomorphous Compounds 

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A new set of relations between structure factors corresponding to an isomorphous pair is established: inequalities, convolution equations, probabilities. These relations may enhance the power of direct methods for the determination of phases.

The aim of this paper is to establish relations between structure factors belonging to pairs of isomorphous crystals. The theory can be applied also to a set of dif-
fraction data obtained for the same crystal by X-ray and neutron diffraction respectively. These relations can be considered as the extension to isomorphous
pairs of the so-called direct methods for solving crystal structures: inequalities, Sayre's equation, theory of probability.

## 1. Inequalities

The unitary structure factors for an isomorphous pair will be denoted by:

$$
\begin{gather*}
U_{\mathbf{H}}=\sum_{j=1}^{N} n_{j} \exp \left(2 \pi i \mathbf{H} \cdot \mathbf{r}_{j}\right),  \tag{1}\\
U_{\mathbf{H}}^{\prime}=\sum_{j=1}^{N} n_{j}^{\prime} \exp \left(2 \pi i \mathbf{H} \cdot \mathbf{r}_{j}\right),  \tag{la}\\
\text { with } \sum_{j} n_{j}=\sum_{j} n_{j}^{\prime}=1, \tag{2}
\end{gather*}
$$

where $n^{\prime}$ is a complex or a real number, $n$ is a real, strictly positive number, $N$ is the number of atoms in the unit cell.
(a) Inequalities arising from Cauchy's inequality

These inequalities are an extension of the Harker \& Kasper (1948) inequalities. We recall first the Cauchy's inequality (3) verified by any pair of sets of complex numbers $a_{j}$ and $b_{j}(j=1, \ldots n)$ :

$$
\begin{equation*}
\left|\sum_{j=1}^{N} a_{j} b_{j}\right|^{2} \leq\left[\sum_{j=1}^{N}\left|a_{j}\right|^{2}\right]\left[\sum_{j=1}^{N}\left|b_{j}\right|^{2}\right] . \tag{3}
\end{equation*}
$$

We now consider a centrosymmetrical isomorphous pair; the unitary structure factors are given by:

$$
\begin{align*}
& U_{\mathbf{H}}=\sum_{j=1}^{N} n_{j} \cos \left(2 \pi \mathbf{H} \cdot \mathbf{r}_{j}\right),  \tag{4}\\
& U_{\mathbf{H}}^{\prime}=\sum_{j=1}^{N} n_{j}^{\prime} \cos \left(2 \pi \mathbf{H} \cdot \mathbf{r}_{j}\right) . \tag{4a}
\end{align*}
$$

Let us apply Cauchy's inequality by setting:
$\ldots a_{j}=n_{j}^{\prime} / n_{j}^{t} \quad \ldots b_{j}=n_{j}^{t} \cos \left(2 \pi \mathbf{H} . \mathbf{r}_{j}\right)$.
We obtain:

$$
\begin{aligned}
\mid \sum_{j}\left(n_{j}^{\prime} / n_{j}^{1 / 2}\right) n_{j}^{1 / 2} & \left.\cos \left(2 \pi \mathbf{H} \cdot \mathbf{r}_{j}\right)\right|^{2} \\
& \leq\left[\sum_{j}\left|n_{j}^{\prime}\right|^{2} / n_{j}\right]\left[\sum_{j} n_{j} \cos ^{2} 2 \pi \mathbf{H} \cdot \mathbf{r}_{j}\right]
\end{aligned}
$$

or

$$
\left|\sum_{j} n_{j}^{\prime} \cos \left(2 \pi \mathbf{H} \cdot \mathbf{r}_{j}\right)\right|^{2} \leq A \sum_{j} n_{j}\left(\frac{1+\cos 4 \pi \mathbf{H} \cdot \mathbf{r}_{j}}{2}\right)
$$

with

$$
\begin{equation*}
A=\sum_{j} \frac{\left.\left|n_{j}^{\prime}\right|\right|^{2}}{n_{j}} . \tag{6}
\end{equation*}
$$

By using (1), we obtain finally:

$$
\begin{equation*}
\left|U_{\mathbf{H}}^{\prime}\right|^{2} \leq A\left(\frac{1}{2}+\frac{1}{2} U_{2 \mathbf{H}}\right) . \tag{7}
\end{equation*}
$$

This inequality enables one to obtain information about the phase of $U_{2 \mathrm{H}}$ by using the knowledge of the modulus of $U_{\mathbf{H}}^{\prime}$. By following the same reasoning we obtain similar inequalities for different space groups. For example, in space group $P 2_{1}$, by setting:

$$
\begin{gathered}
a_{j}=\frac{n_{j}^{\prime}}{n_{j}^{2}} \exp \left(2 \pi i k y_{j}\right), \\
b_{j}=n_{j}^{1 / 2}\left\{\exp \left[2 \pi i\left(h x_{j}+l z_{j}\right)\right]\right. \\
\left.+(-1)^{k} \exp \left[-2 \pi i\left(h x_{j}+l z_{j}\right)\right]\right\}
\end{gathered}
$$

we obtain:

$$
\begin{equation*}
\left|U_{\left.h k\right|^{\prime}}\right|^{2} \leq A\left(\frac{1}{2}+\frac{1}{2}(-1)^{k} U_{2 h \cdot 0 \cdot 21}\right) . \tag{8}
\end{equation*}
$$

## (b) Determinants associated with isomorphous pairs

The inequalities of § (a) may be derived from more general inequalities involving determinants, in the same way as the Harker-Kasper inequalities may be derived from Karle-Hauptman determinants. In order to establish these inequalities, we use the $N$-dimensional Hilbert space, described by Kitaigorodski (1961).* However, the $m$ th vector is now defined differently, as follows:
$p \leq m-1: \quad \mathbf{V}_{p}=\sum_{j}\left[n_{j}^{1 / 2} \exp \left(2 \pi i \mathbf{H}_{p} \cdot \mathbf{r}_{j}\right)\right] \mathrm{e}_{j} ;$
$p=m: \quad \mathbf{W}=\mathbf{V}_{m}=\sum_{j}\left[\frac{n_{j}^{\prime}}{n_{j}^{1 / 2}} \exp \left(2 \pi i \mathbf{H}_{m} \cdot \mathbf{r}_{j}\right)\right] \mathbf{e}_{j}$.
We then form the scalar products:

$$
\begin{aligned}
\left(\mathbf{V}_{p} \cdot \mathbf{V}_{q}\right) & =\sum_{j} n_{j} \exp \left[2 \pi i\left(\mathbf{H}_{p}-\mathbf{H}_{q}\right) \cdot \mathbf{r}_{j}\right]=U_{p q} \\
p, q & =1 \ldots m-1 \\
\left(\mathbf{V}_{p} \cdot \mathbf{V}_{p}\right) & =\sum_{j} n_{j}=1 \\
\left(\mathbf{V}_{p} \cdot \mathbf{W}\right) & =\sum_{j} n_{j}^{\prime} \exp \left[2 \pi i\left(\mathbf{H}_{p}-\mathbf{H}_{m}\right) \cdot \mathbf{r}_{j}\right]=U_{p m}^{\prime} \\
p & =1, \ldots m-1 \\
(\mathbf{W} \cdot \mathbf{W}) & =A \quad(A \text { is given by equation } 6) .
\end{aligned}
$$

With these scalar products we construct the following determinant of order $m$ denoted by $I_{m}$ ( $I$ for 'isomorphous pair').

$$
I_{m}=\left\lvert\, \begin{array}{cccc}
U_{11} & U_{12} & \cdots & U_{1 m}^{\prime} \\
U_{21} & U_{22} & \cdots & U_{2 m}^{\prime} \\
\ldots & \cdots & \cdots & \cdots \\
U_{m 1}^{\prime} & U_{m 2}^{\prime} & \cdots & A
\end{array}\right.
$$

$I_{m}$ is a Gram determinant; therefore:

| for $m \geq N+1$ | $I_{m}=0$, |
| :--- | :--- |
| for $m \leq N$ | $I_{m} \geq 0$. |

## Remark

If both $n_{j}$ and $n_{j}^{\prime}$ are positive, the role of $U_{\mathbf{H}}$ and $U_{\mathbf{H}}^{\prime}$ can be inverted in all equations derived in this paper.

## 2. Algebraic equations

By generalization of the reasoning which leads to the Sayre's (1952) equation, we obtain a similar relation for isomorphous pairs; in this paper we will use the form given by Hughes (1953). We introduce the normalized structure factor $E_{\mathrm{H}}$ defined by

[^0]\[

$$
\begin{gather*}
E_{\mathbf{H}}=\sum_{j=1}^{n} g_{j} \exp \left(2 \pi i \mathbf{H} \cdot \mathbf{r}_{j}\right)  \tag{12}\\
\sum_{j=1}^{N} g_{j}^{2}=\sum_{j=1}^{N} g_{j}^{\prime 2}=1, \quad \mid \overline{\left.E_{\mathbf{H}}\right|^{2}}(\mathbf{H})=\overline{\left|E_{\mathbf{H}}^{\prime}\right|^{2}}(\mathbf{H})=1 \tag{15}
\end{gather*}
$$
\]

The $g_{j}$ 's are related to the $n_{j}$ 's or $f_{j}$ 's (atomic scattering factors) by:

$$
g_{j}=\frac{n_{j}}{\left(\sum_{j} n_{j}^{2}\right)^{1 / 2}}=\frac{f_{j}}{\left(\sum_{j} f_{j}^{2}\right)^{1 / 2}}
$$

and the horizontal bar means 'average value of the expression underneath' when $\mathbf{H}$ sweeps all reciprocal space.

By multiplying $E_{\mathbf{K}}$ and $E_{\mathbf{H}-\mathbf{K}}$, and averaging over $\mathbf{K}$ ( $\mathbf{H}$ being kept constant) we obtain:

$$
\begin{equation*}
{\overline{E_{\mathbf{K}}} \bar{E}_{\mathbf{H}-\mathbf{K}}^{\prime}}^{(\mathbf{K})}=\sum_{j=1}^{N} g_{j} g_{j}^{\prime} \exp \left(2 \pi i \mathbf{H} \cdot \mathbf{r}_{j}\right)=G_{\mathbf{H}} \tag{13}
\end{equation*}
$$

$G_{\mathbf{H}}$ is the Fourier coefficient corresponding to the product of the electron densities $\varrho(\mathbf{r})$ and $\varrho^{\prime}(\mathbf{r})$.

In the general case, $G_{\mathbf{H}}$ is not simply related to $E_{\mathbf{H}}$ or $E_{\mathbf{H}}^{\prime}$; in the same way, the expression

$$
{\overline{E_{\mathrm{K}}} E_{\mathrm{H}-\mathrm{K}}}^{(\mathrm{K})}
$$

is not simply related to $E_{\mathbf{H}}$ unless all atoms are equal.
In our problem, if all atoms in $\varrho(\mathbf{r})$ [not necessarily in $\left.\varrho^{\prime}(\mathbf{r})\right]$ are equal, we have:

$$
\begin{equation*}
G_{\mathrm{H}}=N^{-1 / 2} E_{\mathrm{H}}^{\prime}={\overline{E_{\mathrm{K}}} E_{\mathrm{H}-\mathrm{K}}^{\prime}}_{(\mathrm{K})} \tag{14}
\end{equation*}
$$

Equation (14) suggests that the relationship (in the centrosymmetrical case):

$$
E_{-\mathbf{H}}^{\prime} E_{\mathbf{K}} E_{\mathbf{H}-\mathbf{K}}^{\prime} \sim \text { positive }
$$

is likely to be true even if it is not strictly required by the inequalities. To answer in a quantitative way the question 'How probable is it that (15) holds?' the theory of probability is used.

## 3. Statistical relations

By using the mathematical apparatus described in the paper by Tsoucaris (1970), we obtain the following approximate expression for the probability $P_{+}$that (15) holds:

$$
\begin{equation*}
P_{+}=\frac{1}{2}+\frac{1}{2} \tanh \left[\frac{1}{\sqrt{N}}\left|E_{\mathbf{H}}^{\prime} E_{\mathbf{K}} E_{-\mathbf{H}-\mathrm{K}}\right|\right] . \tag{16}
\end{equation*}
$$

Equation (16) is still approximately true for the general case where the atoms in $\varrho(\mathbf{r})$ are unequal (see Cochran \& Woolfson, 1955).

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# Sesquioxyde de Plomb, $\mathrm{Pb}_{2} \mathbf{O}_{3}$. I. Détermination de la Structure 

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#### Abstract

The structure proposed for lead sesquioxide by Byström (1947) is incorrect. The space group found in the present work is $P 2_{1} / a$ with four $\mathrm{Pb}_{2} \mathrm{O}_{3}$ in the unit cell; the unit-cell dimensions are: $a=7.814$, $b=5.625, c=8.466 \AA($ all $\pm 0.003 \AA) ; \beta=124^{\circ} 48^{\prime} \pm 5^{\prime}$. The structure has been determined by means of X-ray and neutron-diffraction powder data, and refined by a full-matrix least-squares method. The $\mathrm{Pb}^{4+}$ ions are surrounded by six oxygen atoms forming deformed octahedra, with an average $\mathrm{Pb}-\mathrm{O}$ distance of $2 \cdot 18 \AA$. These octahedra share edges and corners to form a two-dimensional framework parallel to (001). These layers are interconnected by $\mathrm{Pb}^{2+}$ ions having sixfold coordination, with $\mathrm{Pb}-\mathrm{O}$ separations ranging from 2.30 to $3.00 \AA$. The three shortest distances $\mathrm{Pb}^{2+} \mathrm{O}^{2-}$ give rise to a $\mathrm{Pb}^{2+}$ coordination reminiscent of the threefold coordination of $\mathrm{Pb}^{2+}$ in $\mathrm{Pb}_{3} \mathrm{O}_{4}$.


## Introduction

Le sesquioxyde de plomb, $\mathrm{Pb}_{2} \mathrm{O}_{3}$, fut préparé pour la première fois par Clark, Schieltz \& Quirke (1937) par synthèse hydrothermale. Gross (1941) proposa pour ce composé une maille monoclinique ayant pour paramètres: $a=7,03, b=5,62, c=3,93 \AA$ et $\beta=82^{\circ}$ avec $Z=2$ et $P 2_{1} / m$ ou $P 2_{1}$ comme groupe spatial. Byström
(1944) adopte les hypothèses de Gross et propose une structure dans laquelle les atomes de plomb occupent deux positions 2(e) du groupe spatial $P 2_{1} / m$; l'atome d'oxygène ayant un facteur de diffusion trop faible par rapport à celui de l'atome de plomb, il lui est impossible de déterminer les positions des atomes d'oxygène, par comparaison des intensités calculées et observées; il propose cependant des positions en tenant compte


[^0]:    * Also Tsoucaris (1970).

